

Electron Density and Doppler RMS Phase Fluctuation in the Inner Corona

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Previous work has developed a self-consistent set of solar wind descriptors for the extended corona ($5 r_{\odot} \leq r \leq 1 \text{ AU}$). In this article, observations of the radial dependence of electron density and RMS phase fluctuation are used to construct a similar and symmetrical set of descriptors for the vastly different regime of the inner corona ($r_{\odot} \leq r \leq 5 r_{\odot}$). The article concludes that the applicability of symmetrical coronal descriptors for both the inner and extended corona argues forcefully for the basic validity of the description.

I. Introduction

In a previous article (Ref. 1), extensive work by J. V. Hollweg (Refs. 2 through 4) was shown to provide a framework in which observations of the radial dependence of phase fluctuation and electron density in the extended corona ($5 r_{\odot} \leq r \leq 1 \text{ AU}$) could be combined with various assumptions and conservation of particle flow to produce an internally consistent set of solar wind descriptors. These are reproduced from Ref. 1 for the extended corona as follows:

where

ϕ = RMS phase fluctuation

N_e = electron density

L_t = transverse fluctuation scale

n = electron density fluctuation

ϵ = ratio of electron density fluctuation to (mean) density

v = radial component of the solar wind velocity

a = signal closest approach distance

r = radial distance

r_{\odot} = solar radius

$$\phi(a) \propto a^{-1.3}$$

$$N_e(r) \propto r^{-2.3}$$

$$L_t(r) \propto r$$

$$n(r) \propto r^{-2.3}$$

$$n/N_e = \epsilon; \epsilon \neq \epsilon(r)$$

$$v(r) \propto r^{0.3}$$

It should be noted that the value for $\phi(a)$ assumes a signal path integration of $(-\infty, \infty)$ about the signal closest approach point; due to the finite limits of the Earth and the spacecraft, the observed radial dependence of actual data will always be somewhat steeper than $a^{-1.3}$.

The inner corona (here to be defined as $r_{\odot} \leq r \leq 5r_{\odot}$) forms a sharply distinct regime ("solar wind acceleration region") from the extended corona in regard to the radial dependence of electron density, solar wind radial velocity, etc., and hence provides an excellent opportunity for validating the extended corona description. It is therefore the purpose of this article to explore the relationship of the various coronal descriptors of the inner corona. Experimental observations of phase fluctuation and electron density in the inner corona will be shown to provide a self-consistent set of coronal descriptors which is symmetrical with those of the extended corona. The fact that the same description prevails in both (vastly distinct) coronal regimes is interpreted as compelling evidence for the validity of the description.

II. Phase Fluctuation Observations in the Inner Corona

The equatorial electron density model derived by Berman and Wackley (Refs. 5 through 7) from Viking Doppler noise and concurrent measurements of dual frequency (S minus X) range is:

$$N_e(r) = \frac{2.39 \times 10^8}{r^6} + \frac{1.67 \times 10^6}{r^{2.30}}, \text{ electrons/cm}^3$$

r = radial distance, solar radii

The doppler noise model ("ISED"; Ref. 5) which is the (signal path integrated) equivalent of the above electron density model is:

$$\text{ISED, Hz} = A_0 \left[\frac{\beta}{(\sin \alpha)^{1.30}} \right] F(\alpha, \beta) + A_1 \left[\frac{1}{(\sin \alpha)^5} \right]$$

$$F(\alpha, \beta) = 1 - 0.05 \left\{ \frac{(\beta - \pi/2 + \alpha)^3 - (\alpha - \pi/2)^3}{\beta} \right\} - 0.00275 \left\{ \frac{(\beta - \pi/2 + \alpha)^5 - (\alpha - \pi/2)^5}{\beta} \right\}$$

where

α = Sun-Earth-probe (SEP) angle, radians

β = Earth-Sun-probe (ESP) angle, radians

$$A_0 = 1.182 \times 10^{-3}$$

$$A_1 = 4.75 \times 10^{-10}$$

For the above models, the electron density terms become equal valued at approximately $r = 4r_{\odot}$, while the corresponding Doppler noise terms become equal valued at approximately $r = 3r_{\odot}$. Doppler noise observations in the region $r < 4r_{\odot}$ are naturally quite sparse; however 20 pass average Doppler values have been accumulated for this inner coronal region from various Solar Conjunctions of Helios 1, Helios 2, and Viking. The method of computing "pass average" values is described in detail in Ref. 8. In deriving the ISEDC model, a three parameter (essentially simultaneous) least squares minimization was performed. The three parameters solved for were the inner corona term coefficient (A_1), and the extended corona term coefficient (A_0) and radial index (-1.30). Figure 1 presents the inner corona data as compared to the inner corona term, the extended corona term, and the combined ISEDC model. The data in the region $r < 4r_{\odot}$ are clearly seen to respond to the (integrated) inner coronal electron density term ($\propto a^{-5}$). Figure 2 presents the correlation between the data and the ISEDC model. Again the correlation is seen to be quite strong. The data in Figures 1 and 2 provide forceful evidence that $\phi(a) \propto a^{-5}$ in the inner corona.

III. Electron Density Observations in the Inner Corona

The radial dependence of electron density in the inner corona has been determined by eclipse photometry methods (van de Hulst, Ref. 9, Saito, Ref. 10, and Blackwell, Ref. 11) and via spacecraft range delays (Muhleman, Mariner 6, Ref. 12, Edenhofer, Helios 2, Ref. 13); in general these experiments have obtained consistent results which indicate a radially dependent electron density in the inner corona of the form $N_e(r) \propto r^{-6}$. Using the relationship derived by Hollweg (Ref. 2), one has for the assumption of a linear transverse scale¹:

$$\phi^2(a) \propto \int_a^{\infty} \frac{1}{(r^6)^2} \frac{r^2 dr}{\sqrt{r^2 - a^2}} \propto (a^{-5})^2$$

¹ The dependence with closest approach distance (a) is obtained by noting the integral:

$$\int_a^{\infty} \frac{r^{\alpha} dr}{\sqrt{r^2 - a^2}}$$

is transformed via the substitution:

$$r = a (\cos x)^{-1}$$

to:

$$a^{\alpha} \int_0^{\pi/2} (\cos x)^{-(1+\alpha)} dx$$

Assuming conservation of particle flow in the corona (Cuperman and Harten, Ref. 14):

$$N_e(r)v(r)r^2 = K$$

and the proportionality between electron density and electron density fluctuation ($\epsilon = n/N_e$), one obtains an inner corona set of self-consistent descriptors symmetrical to those for the extended corona:

$$\phi(a) \propto a^{-5}$$

$$N_e(r) \propto r^{-6}$$

$$L_t(r) \propto r$$

$$n(r) \propto r^{-6}$$

$$n/N_e = \epsilon; \epsilon \neq \epsilon(r)$$

$$v(r) = r^4$$

This set of coronal descriptors is compared to those for the extended corona in Table 1.

IV. Discussion and Conclusions

Previously, observations of the radial dependence of phase fluctuation and electron density in the extended corona were seen to be consistent with the concept of a linear transverse fluctuation scale. The fact that a vastly different regime exists in the inner corona provides a powerful test bed to check the validity of the extended corona description. Observations of doppler noise and electron density in the inner corona, when combined with the assumption of a linear transverse fluctuation scale and conservation of particle flow, are seen to be self-consistent and symmetrical with the extended coronal description. It is hereby concluded that experimental observations in both the inner and extended corona provide compelling evidence for:

- (1) The existence of a linear transverse fluctuation scale.
- (2) The existence of proportionality between RMS phase fluctuation and integrated electron density.

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Table 1. Coronal descriptions for the inner and extended corona

Parameter	Inner corona ($r_{\odot} \leq r \leq 5 r_{\odot}$)	Extended corona ($5 r_{\odot} \leq r \leq 1 \text{ AU}$)
$\phi(a) \propto$	$a^{-5.0}$	$a^{-1.3}$
$N_e(r) \propto$	$r^{-6.0}$	$r^{-2.3}$
$L_t(r) \propto$	r	r
$n(r) \propto$	$r^{-6.0}$	$r^{-2.3}$
$n/N_e =$	$\epsilon; \epsilon \neq \epsilon(r)$	$\epsilon; \epsilon \neq \epsilon(r)$
$\nu(r) \propto$	$r^{4.0}$	$r^{0.3}$

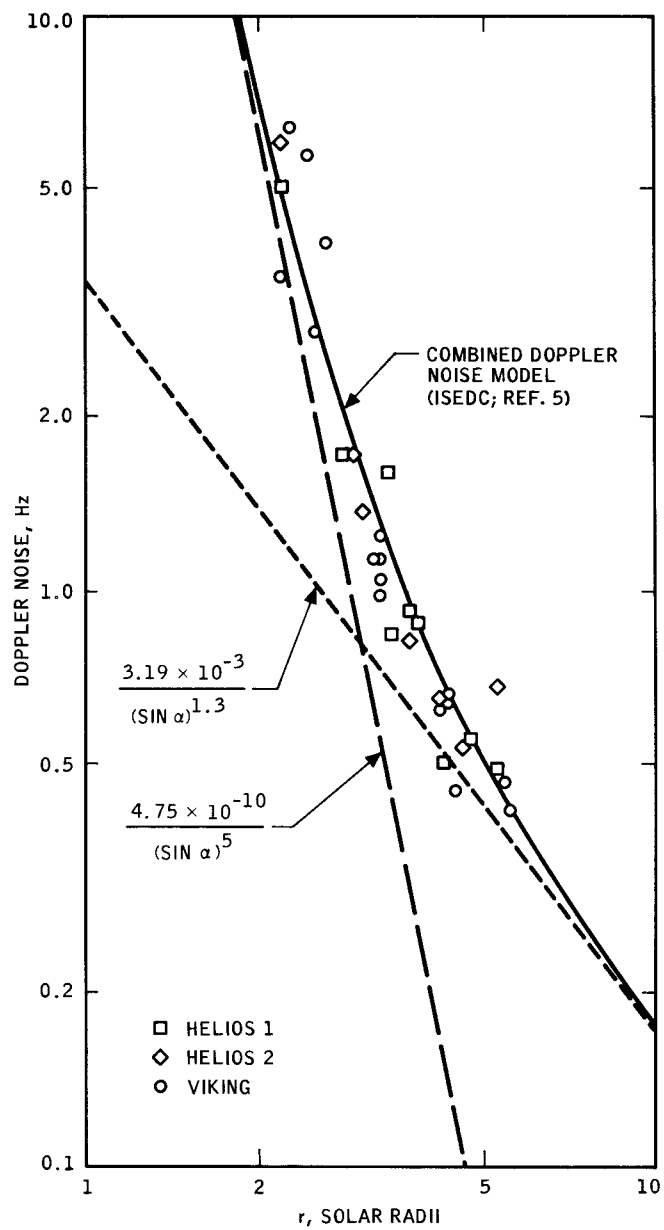


Fig. 1. Doppler noise vs radial distance in the inner corona

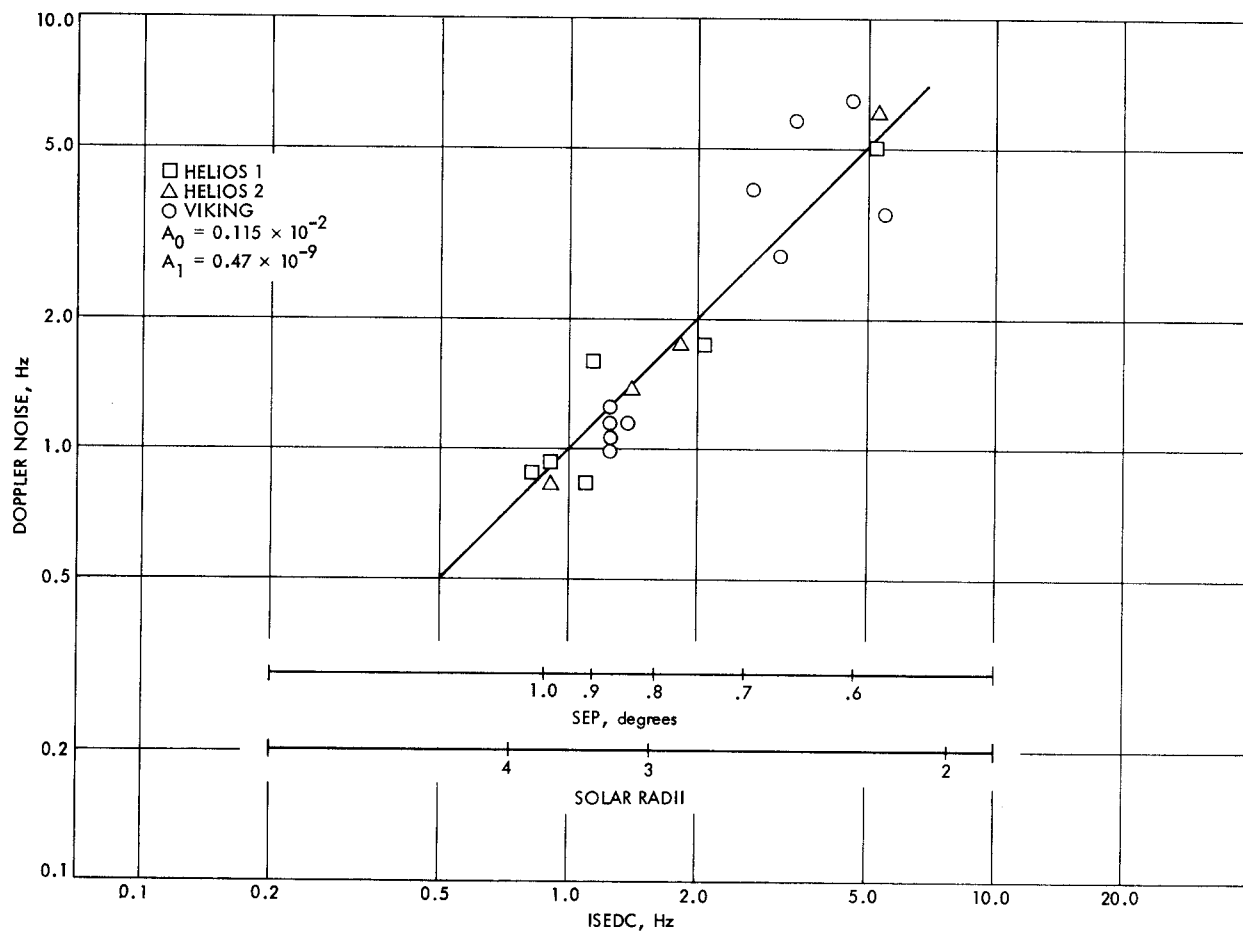


Fig. 2. Correlation of Viking and Helios doppler noise with the ISEDC model between 2 and 4 solar radii